

# Continuous-time Systems

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- LTI systems – recapitulation.
- frequency characteristics  $H(j\omega)$ .
- transfer of a signal through a system with  $H(j\omega)$ .
- Laplace transform.
- Stability and relationship between LT and  $H(j\omega)$ .

## LTI systems – recapitulation

- linearity:  $ax_1(t) + bx_2(t) \longrightarrow ay_1(t) + by_2(t)$ . This property is especially important as we often represent a signal as a sum of complex exponentials.
- time invariance: system's properties do not change over time.
- LTI systems are defined by **impulse response**: for the input  $\delta(t)$  a system outputs  $h(t)$ . What  $h(t)$  looks like for causal systems?
- to any arbitrary input signal  $x(t)$  we compute the output using **convolution**:

$$y(t) = x(t) \star h(t) = \int_{-\infty}^{+\infty} x(\tau)h(t - \tau)d\tau$$

or vice-versa as convolution is communitative. For causal impuls responce we obtain:

$$y(t) = \int_{-\infty}^t x(\tau)h(t - \tau)d\tau$$

## System's response to the input signal $e^{st}$

where  $s$  is an arbitrary complex value:

$$y(t) = H(s)e^{st}, \quad \text{where} \quad H(s) = \int_{-\infty}^{+\infty} h(t)e^{-st} dt.$$

The output is the original signal multiplied by a complex number  $H(s)$ . We are especially interested in the case when  $s = j\omega$ , then the input signal is  $e^{j\omega t}$ .

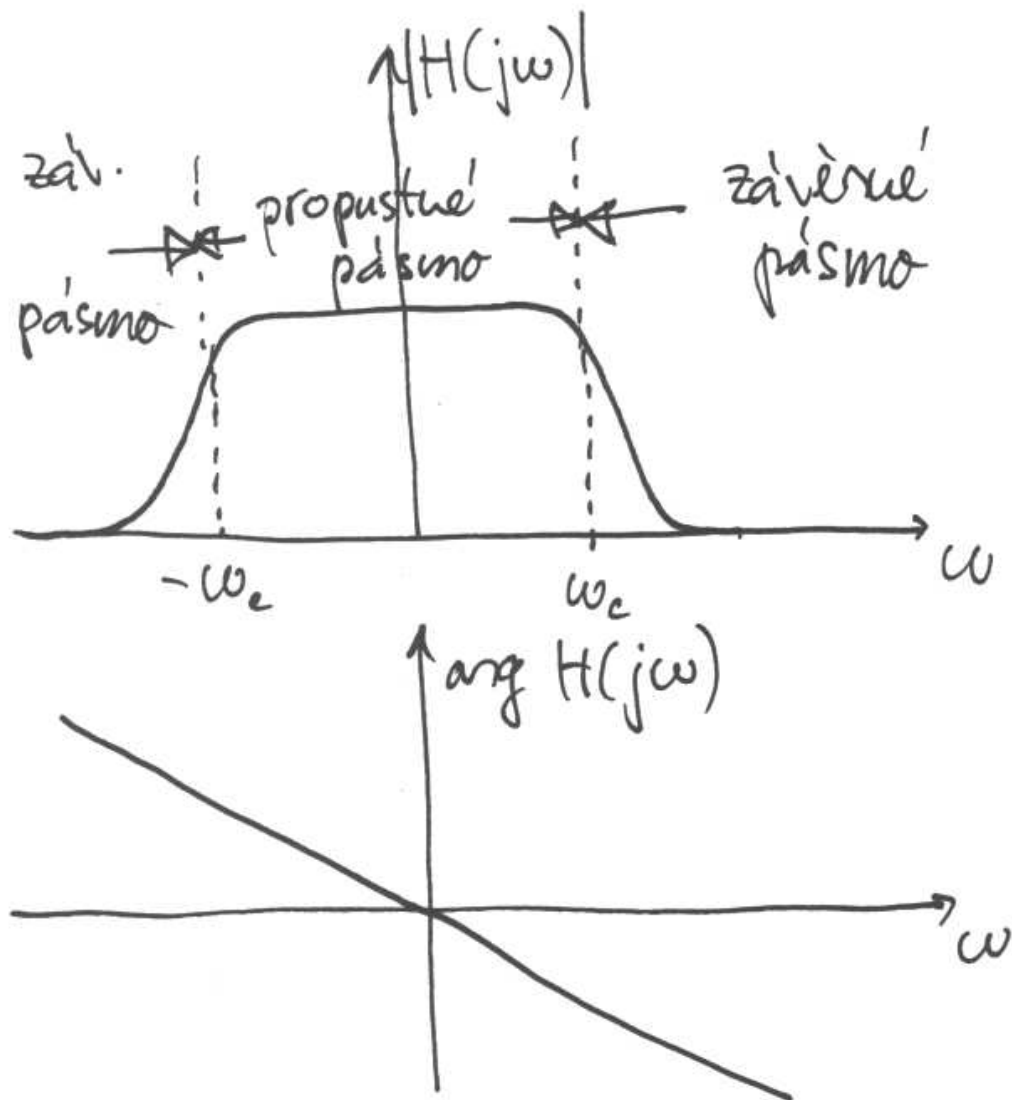
$$y(t) = H(j\omega)e^{j\omega t}, \quad \text{where}$$

$$H(j\omega) = \int_{-\infty}^{+\infty} h(t)e^{-j\omega t} dt.$$

The complex exponential with frequency  $\omega$  is now multiplied by a complex number  $H(j\omega)$ . A value of  $H(j\omega)$  is called **transfer**. We can evaluate  $H(j\omega)$  for any  $\omega$ . The function  $H(j\omega)$  is called **(complex) frequency characteristic**. Transfer function is a Fourier transform of an impulse response:  $H(j\omega) = \mathcal{F}\{h(t)\}$ . As  $h(t) \in \Re$ , then  $H(j\omega)$  has the following property:

$$H(j\omega) = H^*(-j\omega).$$

**Example:**  $H(j\omega)$  of a low pass filter:



## Transfer of a signal through a system with $H(j\omega)$

**Complex exponential**  $x(t) = c_1 e^{j\omega_1 t}$ .

Find a value of  $H(j\omega_1)$ , decompose it into magnitude and argument:

$$y(t) = H(j\omega_1) c_1 e^{j\omega_1 t} = |H(j\omega_1)| |c_1| e^{j(\arg c_1 + \arg H(j\omega_1))} e^{j\omega_1 t}.$$

$\Rightarrow$  only magnitude and argument of  $c_1$  change. The period remains unchanged.

**Cosine function**  $x(t) = C_1 \cos(\omega_1 t + \phi_1)$  can be decomposed to the form:

$x(t) = \frac{C_1}{2} e^{j\phi_1} e^{j\omega_1 t} + \frac{C_1}{2} e^{-j\phi_1} e^{-j\omega_1 t}$ . We work with a **linear system**, that is the exponentials can be processed separately and consequently summed up:

$$y(t) = H(j\omega_1) \frac{C_1}{2} e^{j\phi_1} e^{j\omega_1 t} + H(-j\omega_1) \frac{C_1}{2} e^{-j\phi_1} e^{-j\omega_1 t}.$$

We know that  $H(j\omega_1)$  and  $H(-j\omega_1)$  are complex conjugate, thus  $|H(j\omega_1)| = |H(-j\omega_1)|$  and  $\arg H(-j\omega_1) = -\arg H(j\omega_1)$ .

$$\begin{aligned} y(t) &= |H(j\omega_1)| \frac{C_1}{2} e^{j\phi_1 + j \arg H(j\omega_1)} e^{j\omega_1 t} + |H(j\omega_1)| \frac{C_1}{2} e^{-j\phi_1 - j \arg H(j\omega_1)} e^{-j\omega_1 t} = \\ &= |H(j\omega_1)| C_1 \cos [\omega_1 t + \phi_1 + \arg H(j\omega_1)]. \end{aligned}$$

$\Rightarrow$ the resulting cosine has changed magnitude and phase.

**Example:** Ideal Hi-Fi amplifier amplifies from 0 to 20 kHz:

$$|H(j\omega)| = \begin{cases} 100 & \text{for } 0 \leq |\omega| \leq 40000\pi \\ 1 & \text{for } |\omega| > 40000\pi \end{cases} \quad \arg H(j\omega) = -\frac{\omega}{100000}.$$

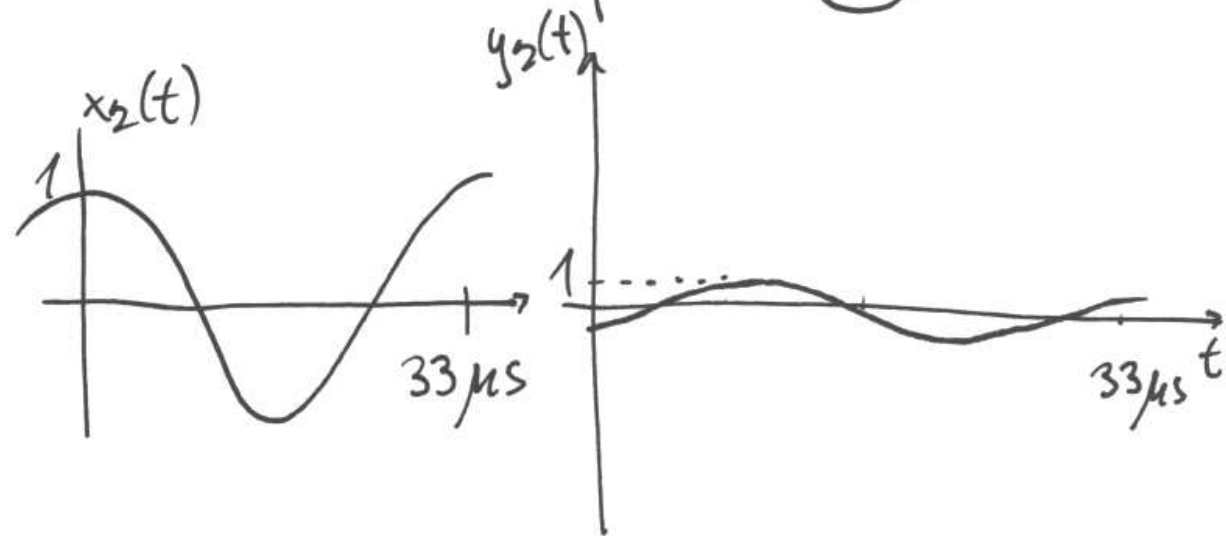
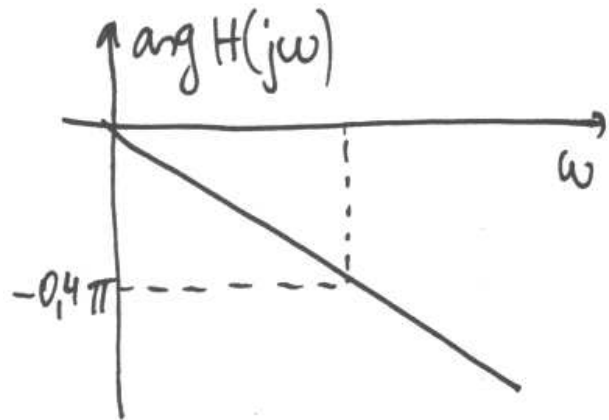
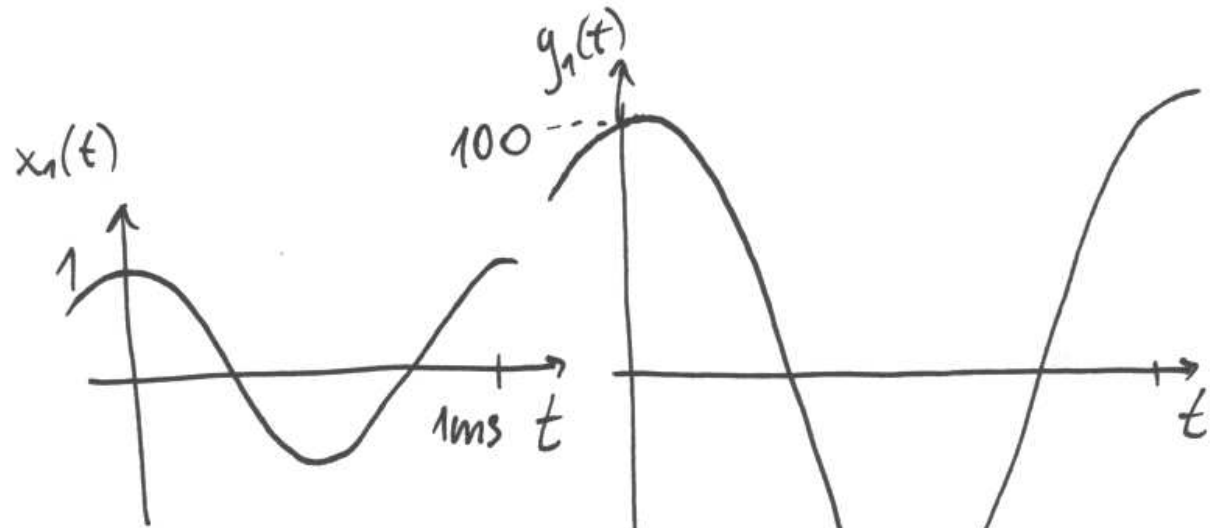
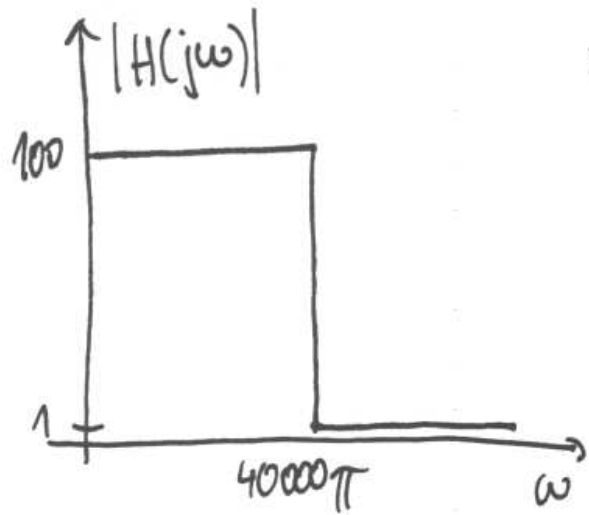
How will it react on a cosine signal with magnitude 1 V and frequencies  $f_1 = 1 \text{ kHz}$  and  $f_2 = 30 \text{ kHz}$  ?

$$x_1(t) = \cos(2000\pi t), \quad \omega_1 = 2000\pi, \quad H(j\omega_1) = 100e^{-j0.02\pi}$$

$$y_1(t) = 100 \cos(2000\pi t - 0.02\pi).$$

$$x_2(t) = \cos(60000\pi t), \quad \omega_2 = 60000\pi, \quad H(j\omega_2) = 1e^{-j0.6\pi}$$

$$y_2(t) = 1 \cos(60000\pi t - 0.6\pi).$$





**Arbitrary periodic signal** can be decomposed into FS:

$$x(t) = \sum_{k=-\infty}^{+\infty} c_k e^{jk\omega_1 t}$$

$H(j\omega)$  will alter every coefficient with respect to the frequency it is tied with:

$$y(t) = \sum_{k=-\infty}^{+\infty} H(jk\omega_1) c_k e^{jk\omega_1 t},$$

Again, we do simple multiplication of the FS coefficients (no convolution).

FS coefficient computation, multiplication and signal restoration can be faster than convolution!

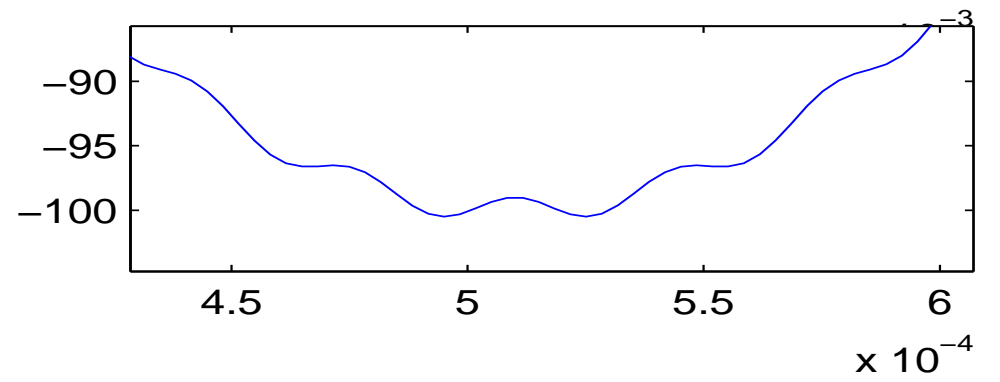
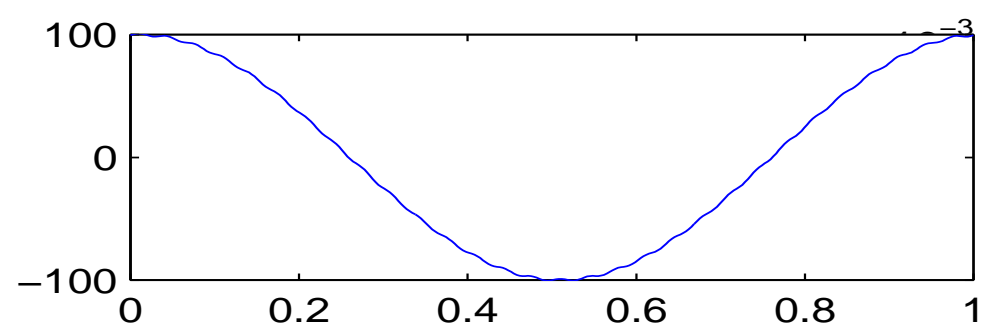
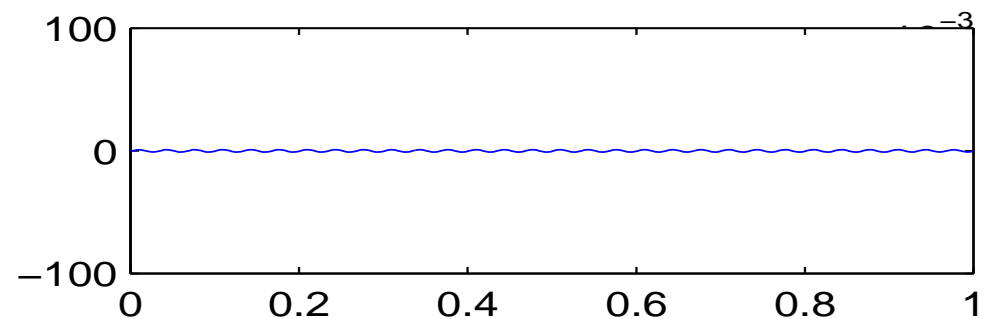
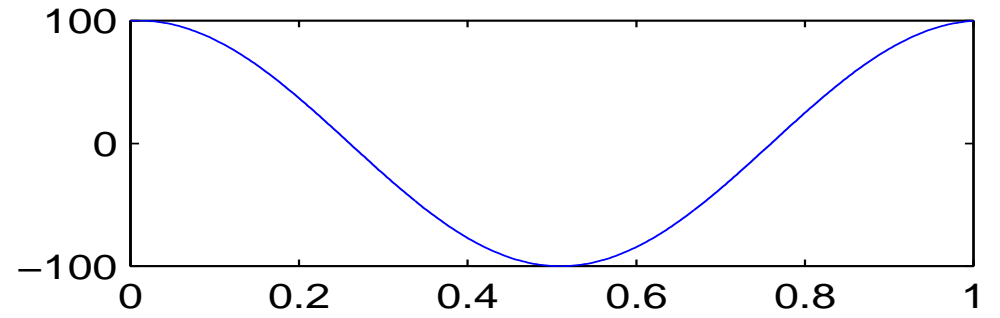
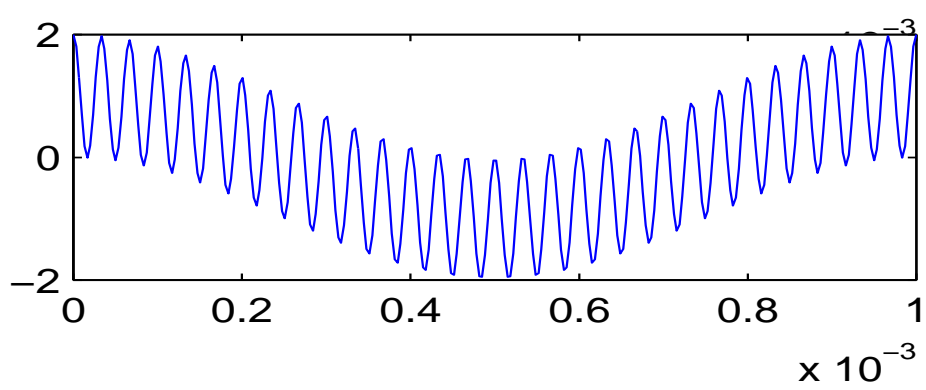
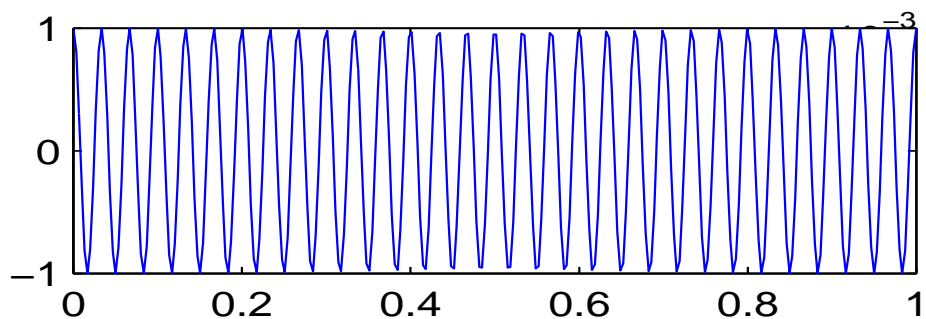
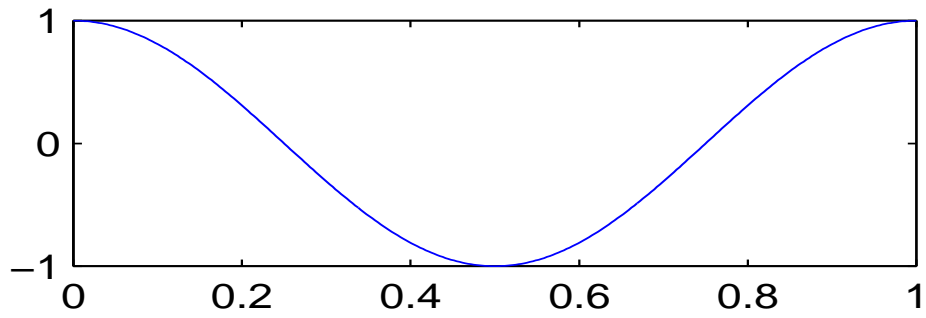
**Example:** A mixture of the signals from the previous example

$$x(t) = \cos(2000\pi t) + \cos(60000\pi t) = c_1 e^{j\omega_1 t} + c_{-1} e^{-j\omega_1 t} + c_{30} e^{j30\omega_1 t} + c_{-30} e^{-j30\omega_1 t},$$

where  $c_1 = c_{-1} = c_{30} = c_{-30} = \frac{1}{2}$ .

After transfer through the amplifier, new coefficients are:

$$c_{1,y} = \frac{1}{2} 100 e^{-j0.02\pi}, \quad c_{-1,y} = \frac{1}{2} 100 e^{j0.02\pi}, \quad c_{30,y} = \frac{1}{2} 1 e^{-j0.6\pi}, \quad c_{-30,y} = \frac{1}{2} 1 e^{j0.6\pi}.$$
$$y(t) = 100 \cos(2000\pi t - 0.02\pi) + 1 \cos(60000\pi t - 0.6\pi).$$



## Transfer of a continuous-time signal through a system with $H(j\omega)$

An arbitrary signal is decomposed to an infinite number of infinitesimally small exponentials using FT:

$$X(j\omega) = \int_{-\infty}^{+\infty} s(t)e^{-j\omega t} dt$$

where  $X(j\omega)$  is a spectral function of input signal and is defined as:

$$X(j\omega) = 2\pi \frac{dc_x}{d\omega}.$$

Likewise, we can define a spectral function of output signal:

$$Y(j\omega) = 2\pi \frac{dc_y}{d\omega}.$$

For some  $\omega_1$  at which we find  $dc_{x,1}$  and  $dc_{y,1}$ , the following holds:

$$dc_{y,1} = H(j\omega_1)dc_{x,1}.$$

And:

$$Y(j\omega_1) \frac{d\omega}{2\pi} = H(j\omega_1) X(j\omega_1) \frac{d\omega}{2\pi}.$$

The above definitions hold for all  $\omega_1$ , thus:

$$Y(j\omega) = H(j\omega)X(j\omega)$$

Can be proved also by:

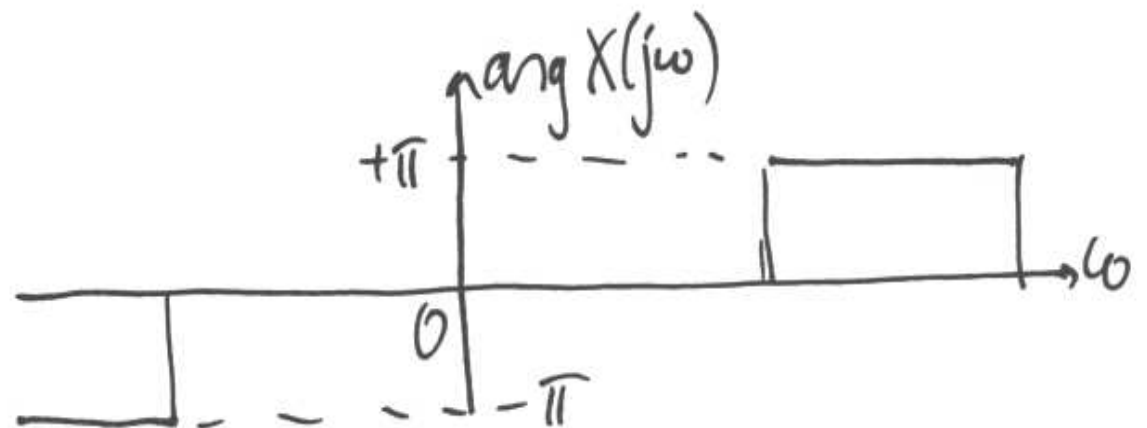
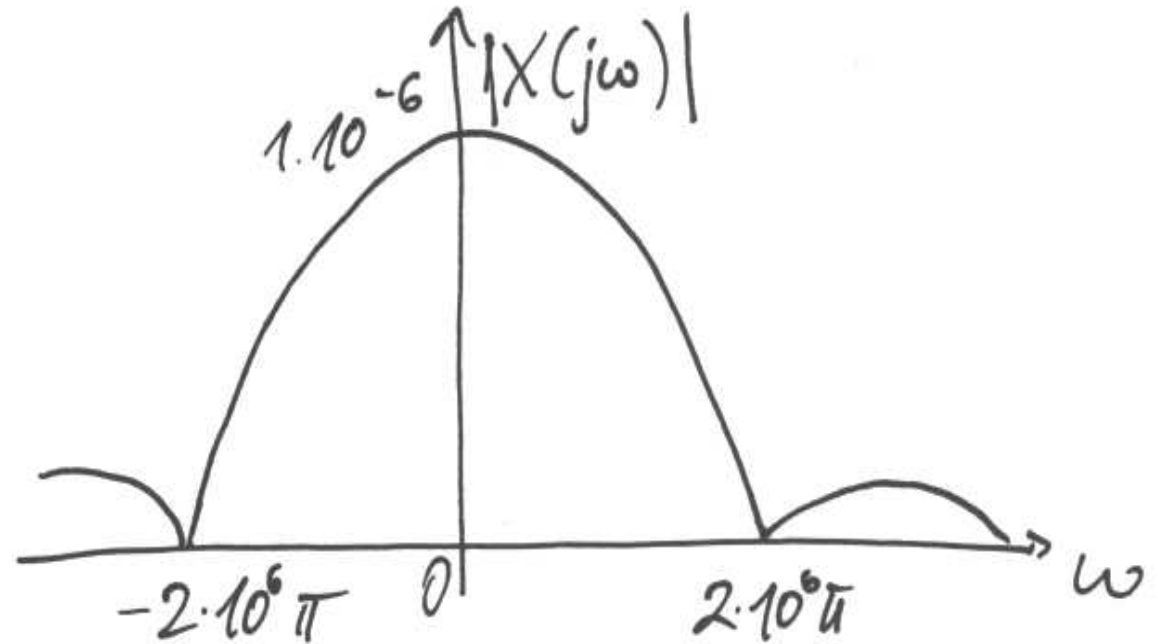
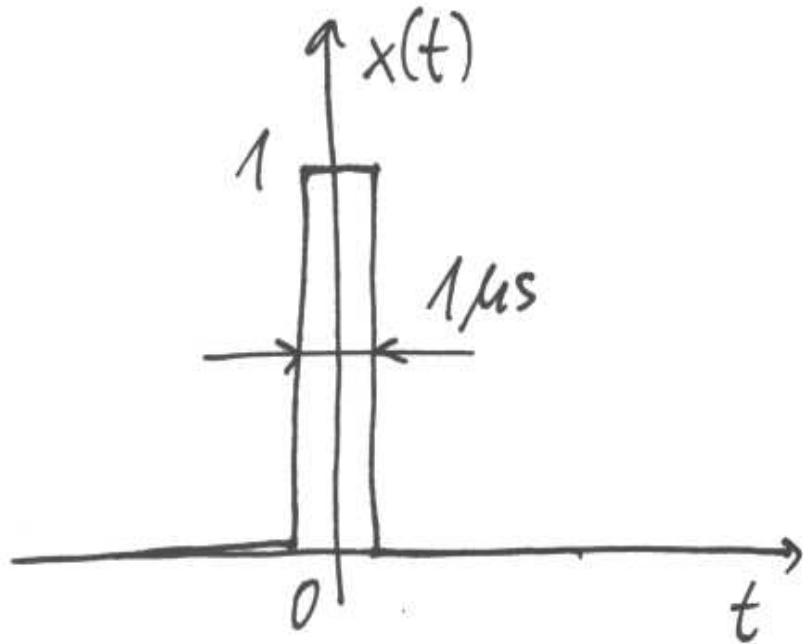
$$Y(j\omega) = \mathcal{F}\{y(t)\} = \int_{-\infty}^{+\infty} \left[ \int_{-\infty}^{+\infty} x(\tau)h(t - \tau)d\tau \right] e^{-j\omega t} dt = \dots$$

**Example:** The input signal into Hi-Fi amplifier is a square signal defined by  $\vartheta = 1 \mu\text{s}$ ,  $D = 1 \text{ V}$ . Calculate the output  $y(t)$

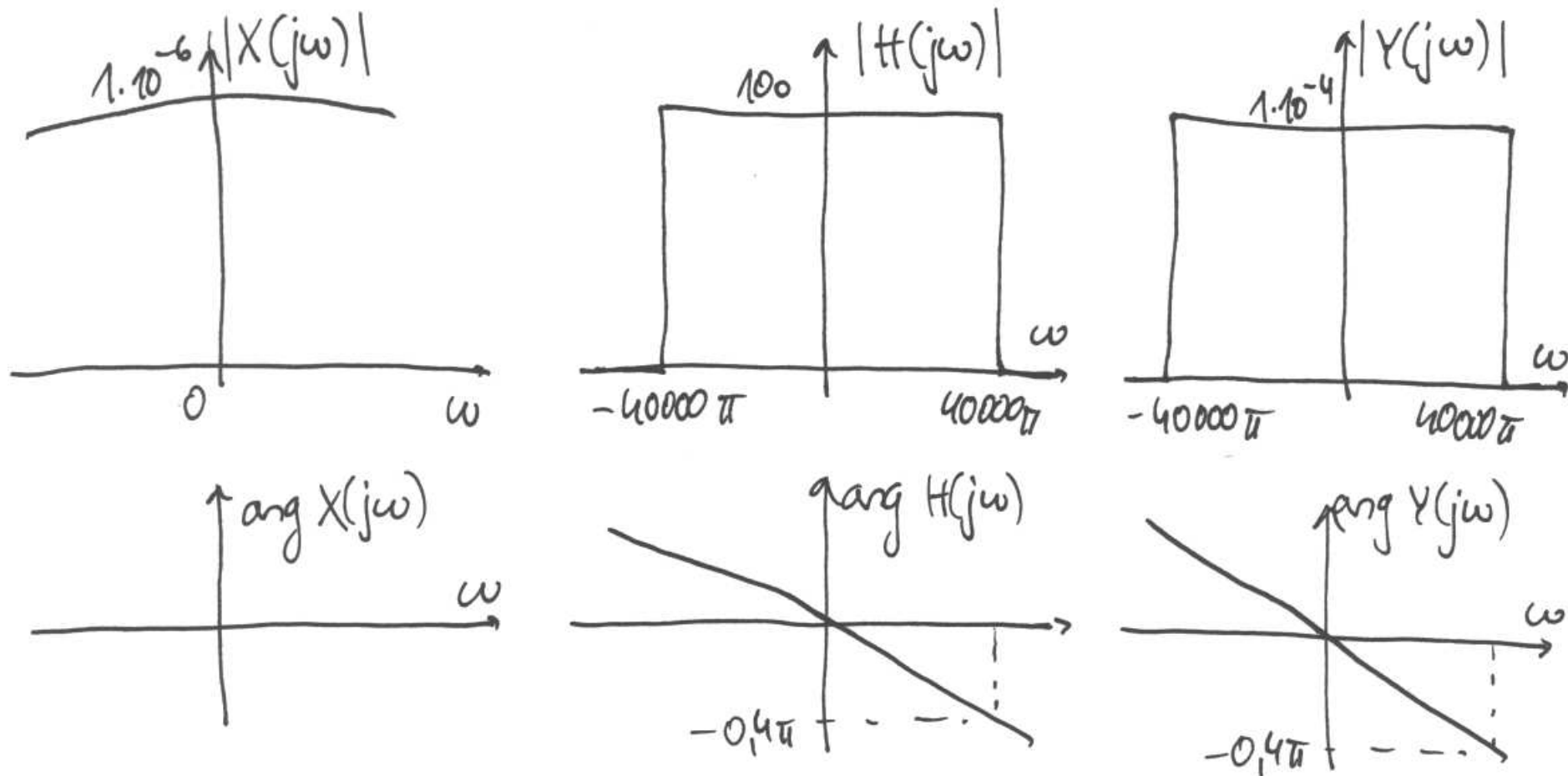
We start off with:

$$x(t) \xrightarrow{\mathcal{F}} X(j\omega) \quad Y(j\omega) = H(j\omega)X(j\omega) \quad Y(j\omega) \xrightarrow{\mathcal{F}^{-1}} y(t).$$

Direct FT:  $X(j\omega) = D\vartheta \text{sinc}\left(\frac{\vartheta}{2}\omega\right) = 1 \times 1 \times 10^{-6} \text{sinc}(0.5 \times 10^{-6}\omega)$ , the function touches the  $\omega$  axis at  $\frac{\vartheta}{2}\omega_a = \pi$ ,  $\omega_a = \frac{\pi}{0.5 \times 10^{-6}} = 2M\pi$ .



The function is multiplied by  $H(j\omega)$ , that is much 'thinner'



The result is a square signal with linear phase. As the last step, an inverse Fourier transform has to be applied.

First, let's assume the phase is zero ("zp" – zero phase). The signal with square shaped spectrum is:

$$y_{zp}(t) = \frac{1}{2\pi} \int_{-\omega_c}^{\omega_c} H e^{+j\omega t} = \frac{H\omega_c}{\pi} \text{sinc}(\omega_c t) = 1 \times 10^{-4} \frac{40000\pi}{\pi} \text{sinc}(40000\pi t) = 4 \text{sinc}(40000\pi t).$$

The function will cross time axis for the first time for:  $40000\pi t = \pi$ ,  $t = \frac{1}{40000} = 25 \mu\text{s}$ .

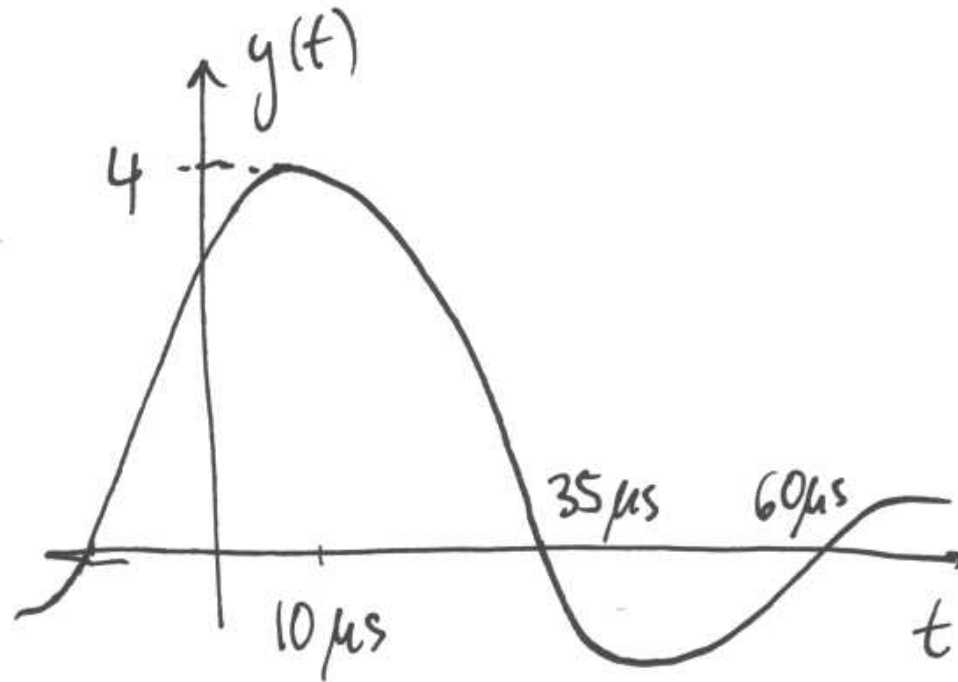
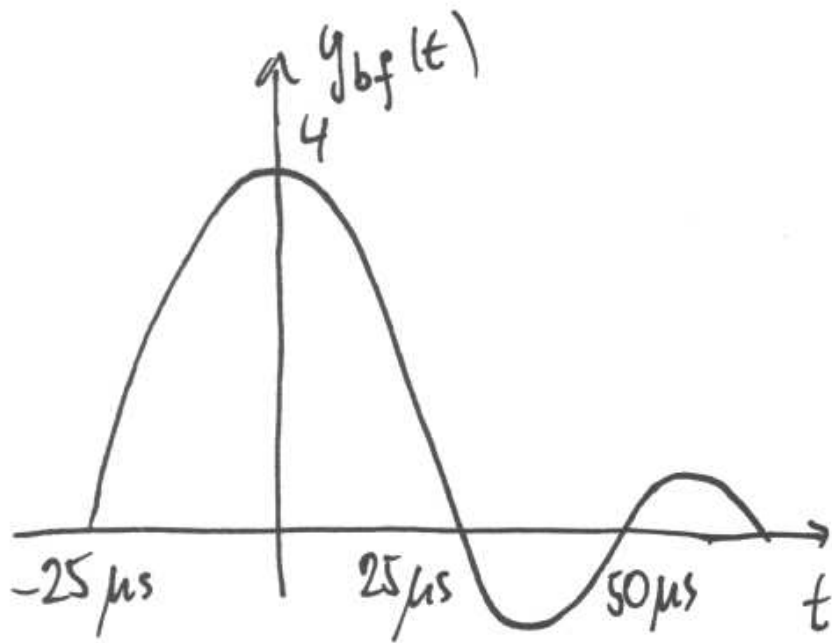
The phase though is linear (not zero-values):  $\phi = -\frac{\omega}{100000} \dots$  What does it remind us?

Shift of the signal in time !

$$y(t) = y_{zp}(t - \tau) \longrightarrow Y(j\omega) = Y_{zp}(j\omega) e^{-j\omega\tau},$$

$$\text{thus } \tau = \frac{1}{100000} = 10 \mu\text{s}.$$



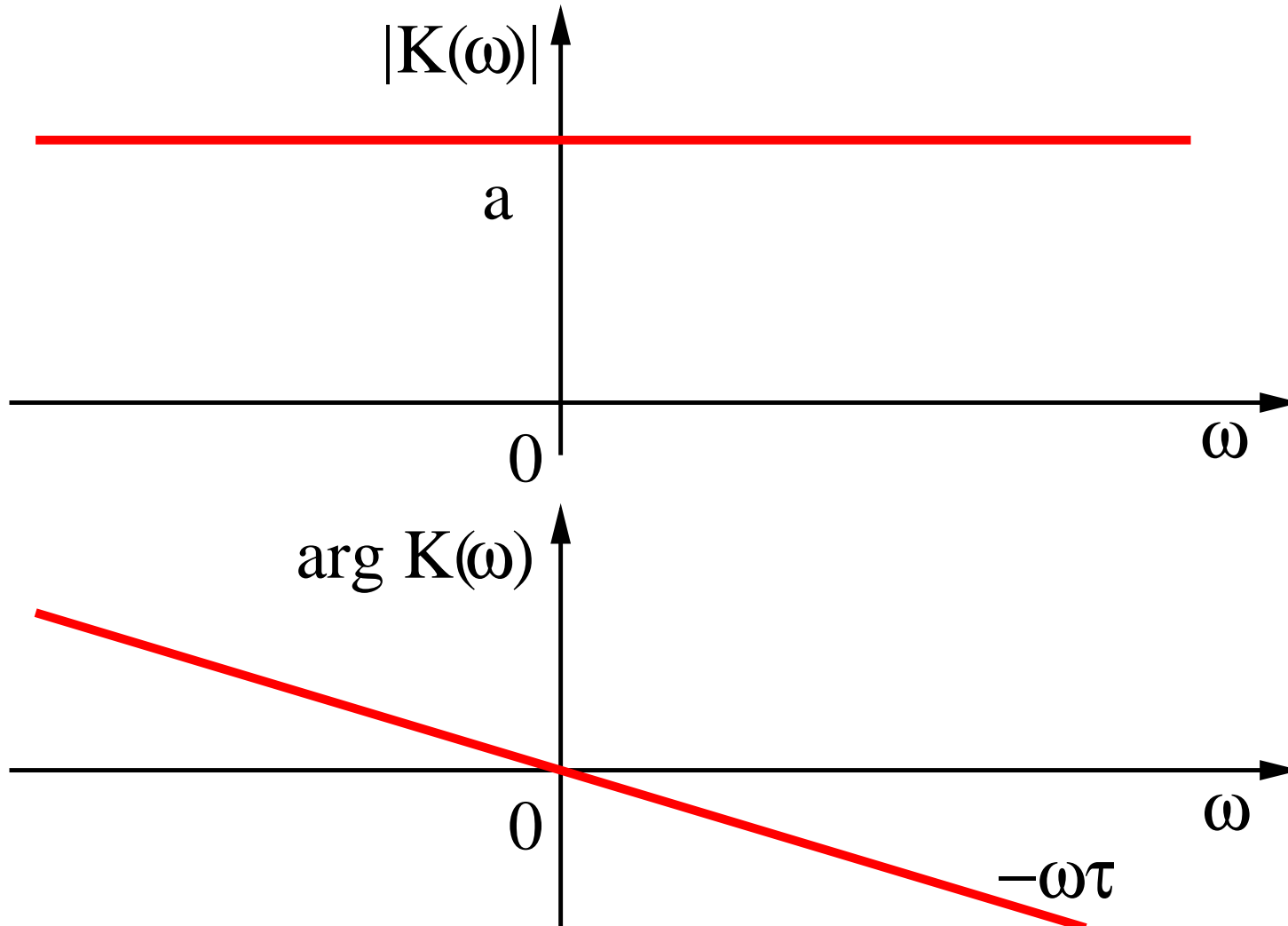


We see that:

- the signal is stretched 50 times.
- low-band filters 'do not like' sharp edged signals in the input.
- impulse response of the ideal low-band filter apparently is function sinc.
- is it really possible for the signal  $y(t)$  begin at time  $-0.5 \mu s$  (before the input  $x(t)$  is actually given to the filter)?

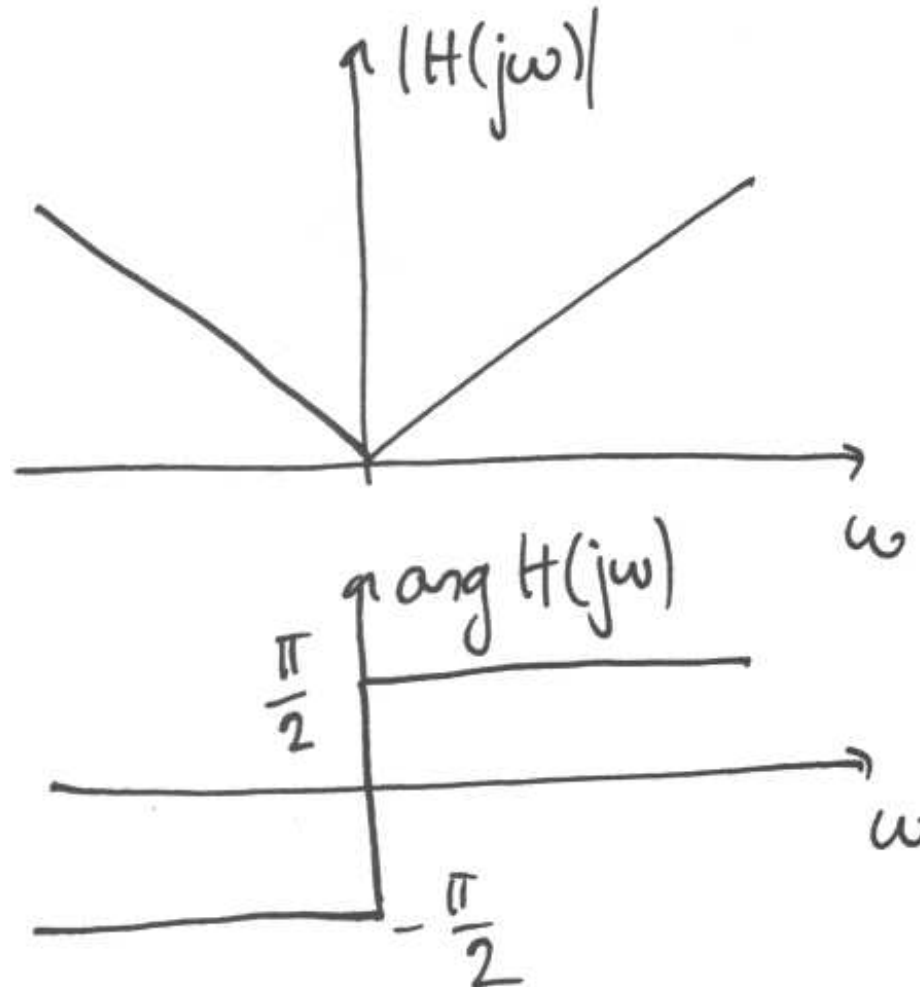
## Ideal amplifier

$$y(t) = a x(t - \tau), \quad Y(\omega) = a X(\omega) e^{-j\omega\tau} \quad H(j\omega) = a e^{-j\omega\tau}$$



## Derivative unit

$y(t) = \frac{dx(t)}{dt}$ . If we set the input  $x(t) = e^{j\omega t}$ , the output is  $y(t) = \frac{de^{j\omega t}}{dt} = e^{j\omega t} j\omega$



$$H(j\omega) = j\omega.$$

## LAPLACE TRANSFORM

We know that given a complex exponential  $e^{st}$  to the input, an LTI system produces the output:

$$y(t) = e^{st} H(s), \quad \text{where} \quad H(s) = \int_{-\infty}^{+\infty} h(\tau) e^{-s\tau} d\tau.$$

So far we were interested in the case when  $s = j\omega$ , now we will be interested in the entire complex plane “ $s$ ”. **Laplace transform:**

$$X(s) = \int_{-\infty}^{+\infty} x(t) e^{-st} dt,$$

where  $s = \sigma + j\omega$  is a complex variable.  $X(s)$  is a complex function over complex plane called **projection** and is denoted as  $\mathcal{L}$ ,  $x(t) \xrightarrow{\mathcal{L}} X(s)$ . Note, that:

$$X(s)|_{s=j\omega} = \mathcal{F}\{x(t)\}.$$

What do we need LT for? We are not interested in the transformed signals but rather the behaviour and stability of the systems defined by differential equations (LTI systems).

## Fundamental properties of LT

- Convolution in time domain  $x_1(t) \star x_2(t) \longrightarrow X_1(s)X_2(s)$ , corresponds to multiplication in frequency domain

$$x(t) \longrightarrow X(s), \quad h(t) \longrightarrow H(s)$$

$$y(t) = x(t) \star h(t) \longrightarrow Y(s) = X(s)H(s).$$

$H(s)$  is called **transfer** function or **system** function.

- Note the derivation  $\frac{dx(t)}{dt} \longrightarrow sX(s)$ .

## Systems defined by differential equations

$$\sum_{k=0}^N a_k \frac{d^k y(t)}{dt^k} = \sum_{k=0}^M b_k \frac{d^k x(t)}{dt^k}$$

We are interested in the transfer function that is telling us what's the system's reaction on the input:

$$\left( \sum_{k=0}^N a_k s^k \right) Y(s) = \left( \sum_{k=0}^M b_k s^k \right) X(s),$$

$$H(s) = \frac{\sum_{k=0}^M b_k s^k}{\sum_{k=0}^N a_k s^k},$$

From the above notation, we can easily get to the frequency characteristics:

$$H(j\omega) = H(s)|_{s=j\omega}$$

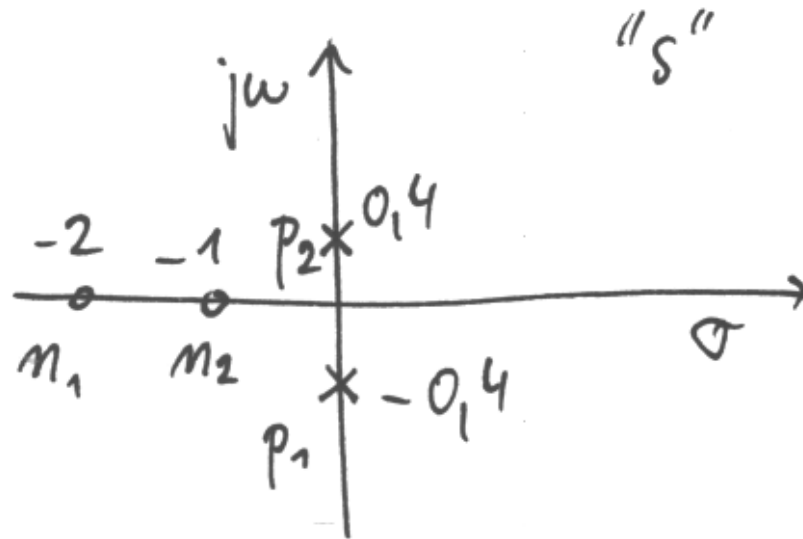
the numerator and the denominator are composed of **polynomials**. Polynomials are defined either by coefficients or **roots**.

$$H(s) = \frac{b_M \prod_{k=1}^M (s - n_k)}{a_N \prod_{k=1}^N (s - p_k)}.$$

The roots are the values of  $s$  for which the polynomial equals to zero. The roots of the numerator  $n_k$  are called **zero points** or **zeros**. The roots of the denominator  $p_k$  are called **poles**.

**Example:**  $H(s) = \frac{s^2+3s+1}{s^2+0.16} = \frac{(s+2)(s+1)}{(s+j0.4)(s-j0.4)}$

$n_1 = -2, n_2 = -1, p_1 = -j0.4, p_2 = +j0.4$



From zeros and poles, we can also estimate graphical representation of  $H(j\omega)$ .



## Stability of causal systems

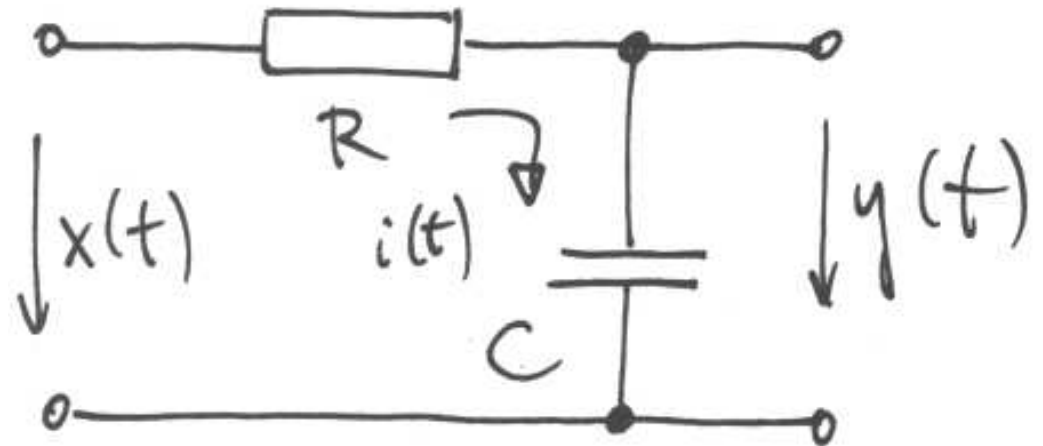
Causal system is stable when all poles lie in the left half-plane of the complex plane:  
 $\Re\{p_k\} < 0$ .

**Example:** what is the capacitor tention  $y(t)$  with respect to the source tention  $x(t)$ ?

$$i(t) = \frac{x(t) - y(t)}{R}$$

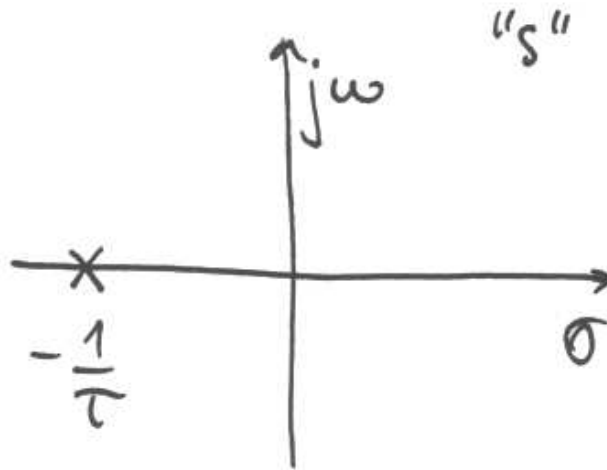
$$\text{current on the capacitor : } i(t) = C \frac{dy(t)}{dt}$$

$$\text{it gives : } RC \frac{dy(t)}{dt} + y(t) = x(t)$$



$\tau = RC$  is so called **time constant** of the system. The coefficients of the differential equation are :  $a_1 = \tau$ ,  $a_0 = 1$ ,  $b_0 = 1$ .

The transfer function has no zeros and one pole for  $(s\tau + 1) = 0$ , thus  $p_1 = -\frac{1}{\tau}$ .



⇒ the system is **stable**.

Zero-pole notation:

$$H(j\omega) = \frac{b_M \prod_{k=1}^M (j\omega - n_k)}{a_N \prod_{k=1}^N (j\omega - p_k)}.$$

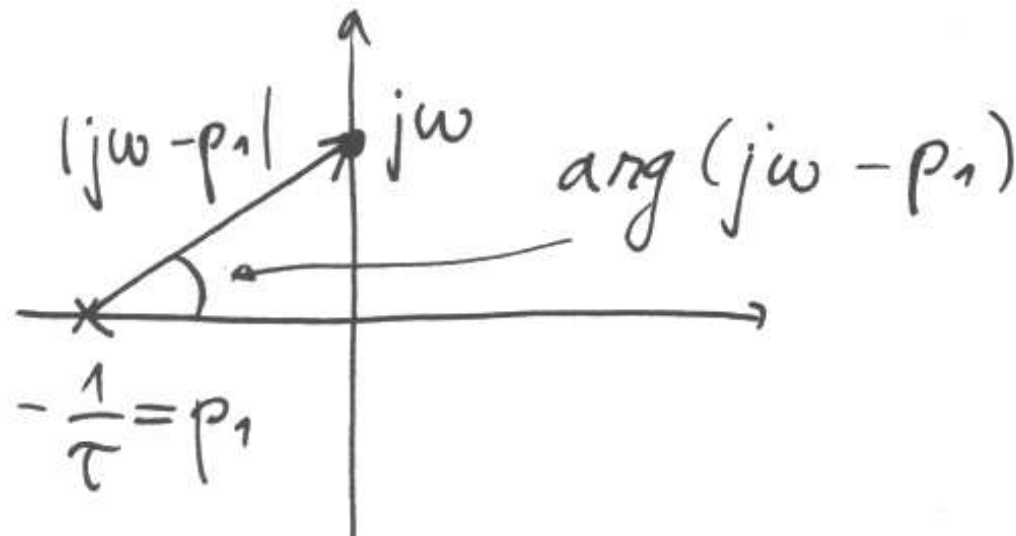
For a given  $j\omega$  each braces is a complex number, that can be represented as a vector.

To obtain the value of the frequency characteristic for the frequency  $\omega$  we have to

- multiply together moduls of all the numbers from the numinator and sum up argumets.
- divide moduls of all the numbers from the denuminator and subtract arguments.

Back to the example...

$$H(s) = \frac{1}{\tau} \frac{1}{s - (-\frac{1}{\tau})}.$$



- we start off with  $j\omega = 0$ :  $|H(0)| = 1$ ,  $\arg H(0) = 0$ .
- increasing  $\omega$ : then  $|s - (-\frac{1}{\tau})|$  increases, but is in the denominator. Thus the fraction decreases. The angle increases and is used with the inverted sign to get the final result.
- $j\omega = \infty$ :  $|H(j\infty)| = 0$ ,  $\arg H(0) = -\frac{\pi}{2}$ .

**Example:**  $R = 1\text{ k}\Omega$ ,  $C = 1\text{ }\mu\text{F}$ ,  $\tau = RC = 1\text{ ms}$ .

